



**ELECTRON-ION RECOMBINATION,
PHOTOIONIZATION AND
DIELECTRONIC SATELLITE LINES
OF Ca XVIII AND Ca XIX USING
UNIFIED METHOD**

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Photoionization and Electron-Ion Recombination



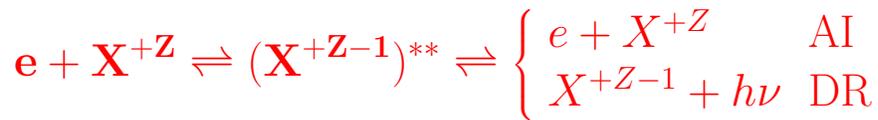
- These bound-free processes are inverse to each other
- Give rise to dielectronic satellite lines (DES)

These may proceed as:

i) Photoionization (PI) & Radiative Recombination (RR)



ii) Autoionization (AI) & Dielectronic Recombination (DR)



- *Doubly excited autoionizing state* $[(\text{X}^{+Z-1})^{**}]$ *introduces resonances*
- RR and DR are inseparable in nature

Theory: "UNIFIED METHOD" - (i) includes RR and DR, (ii) provides total electron-ion recombination rates self-consistent to photoionization cross sections, (iii) dielectronic satellite (DES) line strengths and profiles (*new extension*)

- Computational framework - close coupling approximation and R-matrix method (Nahar & Pradhan, PRL 1992, PRA 1994, PRA 2006)

DIELECTRONIC SATELLITE (DES) LINES

- **DES lines** - common in (e+He-like ion) collision spectra forming 3-electron Li-like ion and are used as diagnostics
- **DES lines** - form below the excitation threshold of the core. Ex. KLL ($1s2l2l'$) satellite lines are below the core excitation $1s^2(^1S_0) \rightarrow 1s2p(^1P_1^o)$ of w-line
- **DES lines** - produced by radiation damping of doubly excited autoionizing states in dielectronic recombination.

Ex. KLL ($1s2l2l'$) complex introduces **22 DES lines** (denoted as a,b,...,v - Gabriel 1972) as autoionizing states $1s2l2l'$ ($1s2p2p$, $1s2s^2$, $1s2p^2$) $\rightarrow 1s^22s(^2S_{1/2})$, $1s^22p(^2P_{1/2,3/2}^o)$ bound states

Theoretical Treatment of DES lines

The earlier treatments

Based on isolated resonance approximation (initiated by Gabriel 1972) where the recombination rate coefficient (Bates & Dalgarno 1962) is given by

$$\alpha_R = a_o^3 \frac{g_i}{2g_f} \left[\frac{4\pi}{T} \right] e^{-\frac{\epsilon}{kT}} \frac{A_r A_a}{\sum_m A_a(m) + \sum_n A_r(n)} \quad (1)$$

ϵ is the DES energy

- Satellite lines are obtained at single energy points, ϵ .

The present treatment

Employs detailed photoionization cross sections for recombination cross sections

$$\sigma_{RC} = \sigma_{PI} \frac{g_i}{g_j} \frac{h^2 \omega^2}{4\pi^2 m^2 c^2 v^2}$$

where σ_{PI} includes autoionizing resonances

- Obtains intensity $I_s = \alpha_s \frac{n_i}{n_e}$, or intensity ratio $\frac{I_s}{I_w} = \frac{\alpha_s}{\alpha_w}$, and resonant line strength $S = \int_{\epsilon_i}^{\epsilon_f} \sigma_{RC} d\epsilon$ from the unified σ_{RC} that includes both RR background and DR
- Provides energy profile of the satellite lines instead of a single point line
- Generates entire DES spectrum naturally

THEORY: Close-coupling (CC) R-matrix method

Total wavefunction expansion in CC approximation:

$$\Psi_{\mathbf{E}}(\mathbf{e} + \mathbf{ion}) = \mathbf{A} \sum_{\mathbf{i}}^{\mathbf{N}} \chi_{\mathbf{i}}(\mathbf{ion}) \theta_{\mathbf{i}} + \sum_{\mathbf{j}} \mathbf{c}_{\mathbf{j}} \Phi_{\mathbf{j}}(\mathbf{e} + \mathbf{ion})$$

$\chi_{\mathbf{i}} \rightarrow$ target/core ion wavefunction,

$\Phi_{\mathbf{j}} \rightarrow$ correlation functions of (e+ion) system

$\theta_{\mathbf{i}} \rightarrow$ electron (bound or continuum) wavefunction

• The complex resonant structures in collisional and radiative processes are included through channel couplings.

The Iron Project - The relativistic Hamiltonian in Breit-Pauli R-matrix (**BPRM**) approximation:

$$\mathbf{H}_{\mathbf{N}+1}^{\mathbf{BP}} = \mathbf{H}_{\mathbf{N}+1}^{\mathbf{NR}} + \mathbf{H}_{\mathbf{N}+1}^{\mathbf{mass}} + \mathbf{H}_{\mathbf{N}+1}^{\mathbf{Dar}} + \mathbf{H}_{\mathbf{N}+1}^{\mathbf{so}}$$

where the non-relativistic Hamiltonian is

$$\mathbf{H}_{\mathbf{N}+1}^{\mathbf{NR}} = \sum_{\mathbf{i}=1}^{\mathbf{N}+1} \left\{ -\nabla_{\mathbf{i}}^2 - \frac{\mathbf{2Z}}{\mathbf{r}_{\mathbf{i}}} + \sum_{\mathbf{j}>\mathbf{i}}^{\mathbf{N}+1} \frac{\mathbf{2}}{\mathbf{r}_{\mathbf{ij}}} \right\}.$$

Mass correction term $\rightarrow \mathbf{H}_{\mathbf{N}+1}^{\mathbf{mass}} = -\frac{\alpha^2}{4} \sum_{\mathbf{i}} \mathbf{p}_{\mathbf{i}}^4$, Darwin term $\rightarrow \mathbf{H}_{\mathbf{N}+1}^{\mathbf{Dar}} = \frac{\mathbf{Z}\alpha^2}{4} \sum_{\mathbf{i}} \nabla^2 \left(\frac{1}{\mathbf{r}_{\mathbf{i}}} \right)$, Spin-orbit interaction term $\rightarrow \mathbf{H}_{\mathbf{N}+1}^{\mathbf{so}} = \mathbf{Z}\alpha^2 \sum_{\mathbf{i}} \frac{1}{\mathbf{r}_{\mathbf{i}}^3} \mathbf{l}_{\mathbf{i}} \cdot \mathbf{s}_{\mathbf{i}}$

Solve Schrodinger equation: $\mathbf{H}_{\mathbf{N}+1}^{\mathbf{BP}} \Psi_{\mathbf{E}}(\mathbf{e} + \mathbf{ion}) = \mathbf{E} \Psi_{\mathbf{E}}(\mathbf{e} + \mathbf{ion})$

The channels introduce a set of coupled equations which are solved by R-matrix method

- $\mathbf{E} < 0 \rightarrow$ Bound (e+ion) states $\Psi_{\mathbf{B}}$
- $\mathbf{E} \geq 0 \rightarrow$ Continuum states $\Psi_{\mathbf{F}}$

Photoionization/Recombination Transition Matrix elements:

$$\langle \Psi_B || \mathbf{D} || \Psi_F \rangle$$

$\mathbf{D} = \sum_n r_n \rightarrow$ dipole operator $n =$ number of electrons

The generalized line strength (\mathbf{S})

$$\mathbf{S} = | \langle \Psi_j || \mathbf{D} || \Psi_i \rangle |^2 = \left| \langle \Psi_f | \sum_{j=1}^{N+1} \mathbf{r}_j | \Psi_i \rangle \right|^2,$$

The photoionization cross section is

$$\sigma_{\text{PI}} = \frac{4\pi}{3c} \frac{1}{g_i} \omega \mathbf{S},$$

Recombination cross section, σ_{RC} , is related to σ_{PI} as,

$$\sigma_{\text{RC}} = \sigma_{\text{PI}} \frac{g_i}{g_j} \frac{h^2 \omega^2}{4\pi^2 m^2 c^2 v^2}.$$

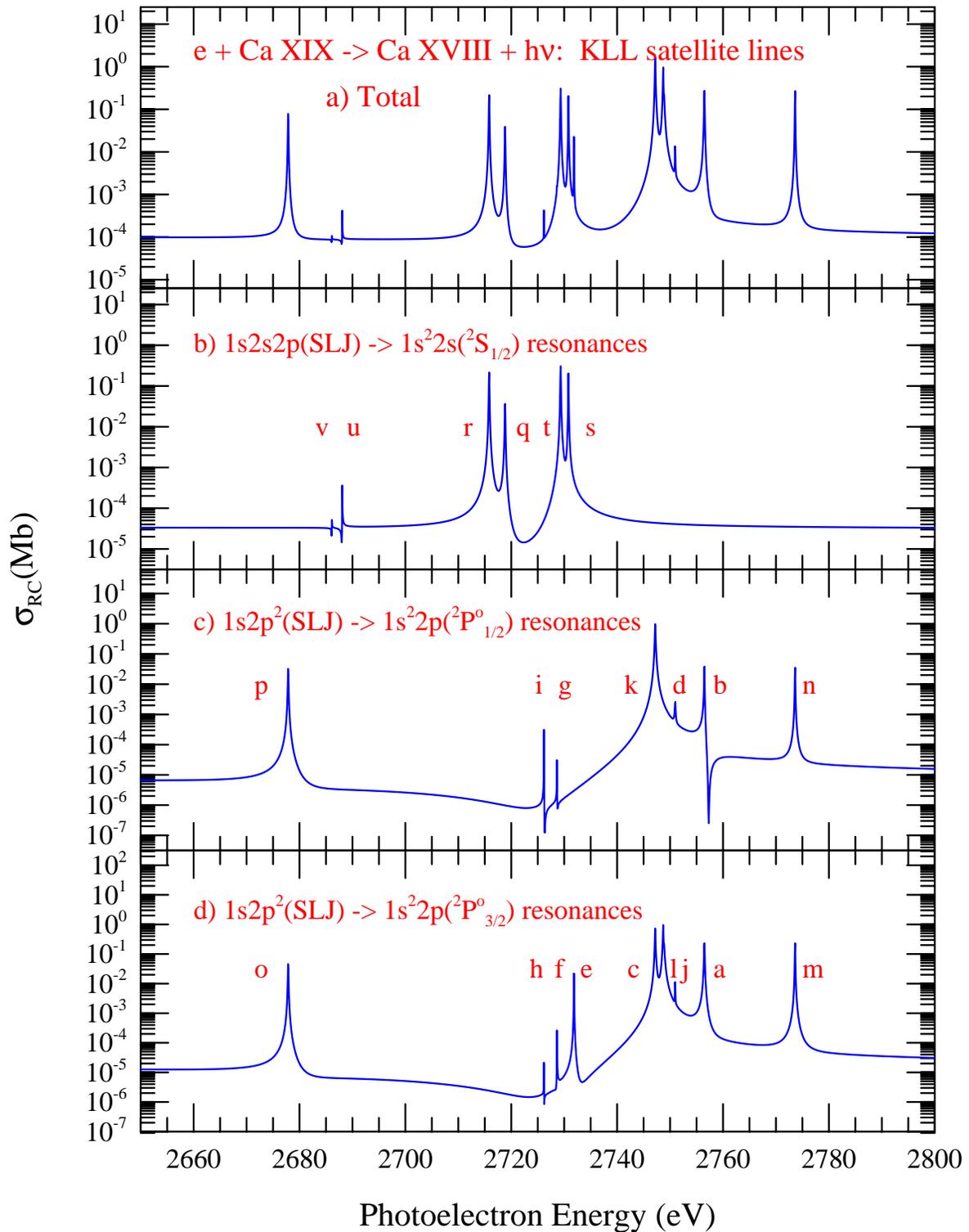
The recombination rate coefficient, α_{RC} , is obtained as

$$\alpha_{\text{RC}}(\mathbf{T}) = \int_0^\infty v f(v) \sigma_{\text{RC}} dv,$$

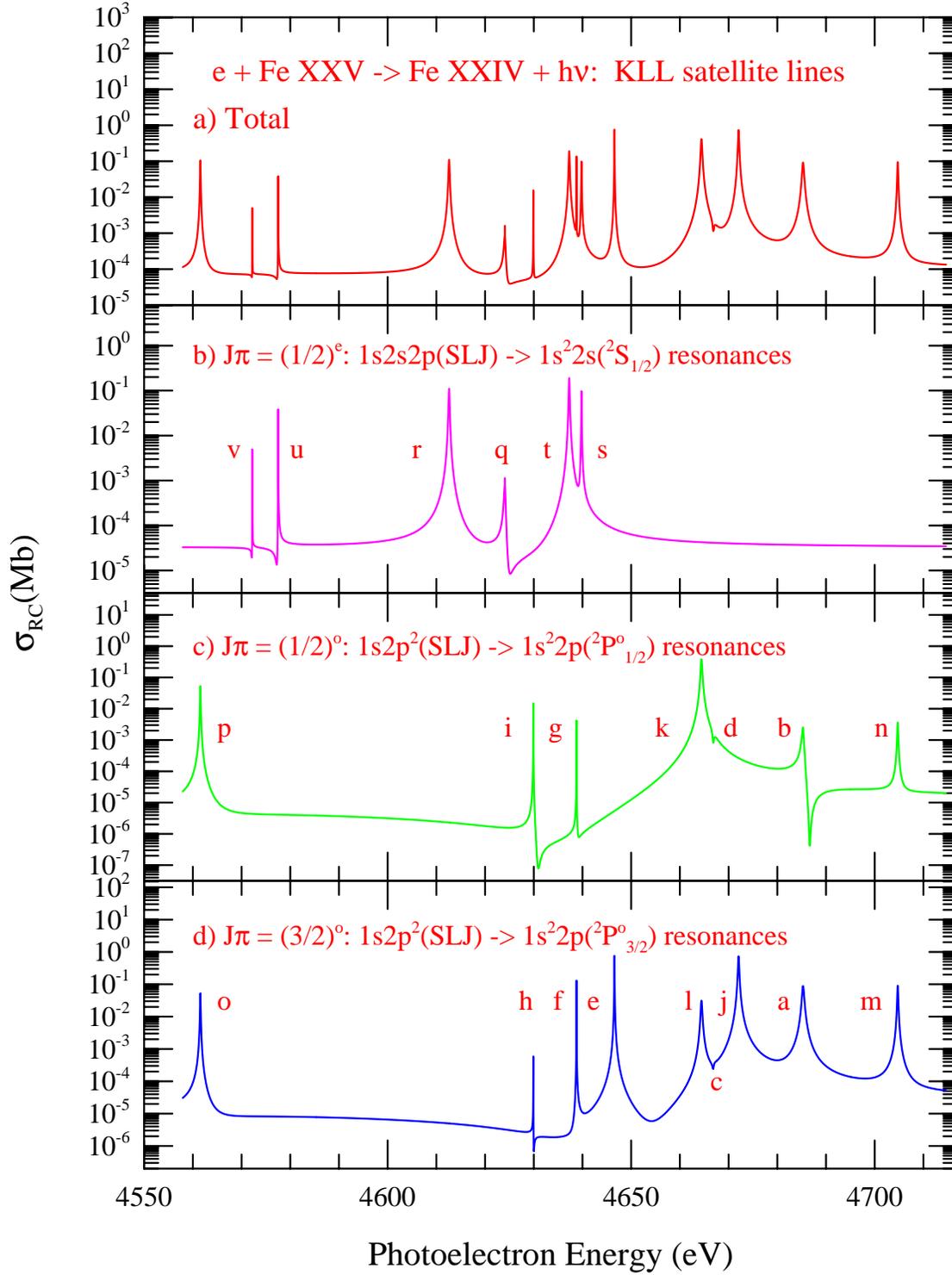
$f(v) =$ Maxwellian velocity distribution function

Total $\alpha_{\text{RC}} \rightarrow$ Contributions from infinite number of recombined states

Dielectronic Satellites (DES) from the Unified Method

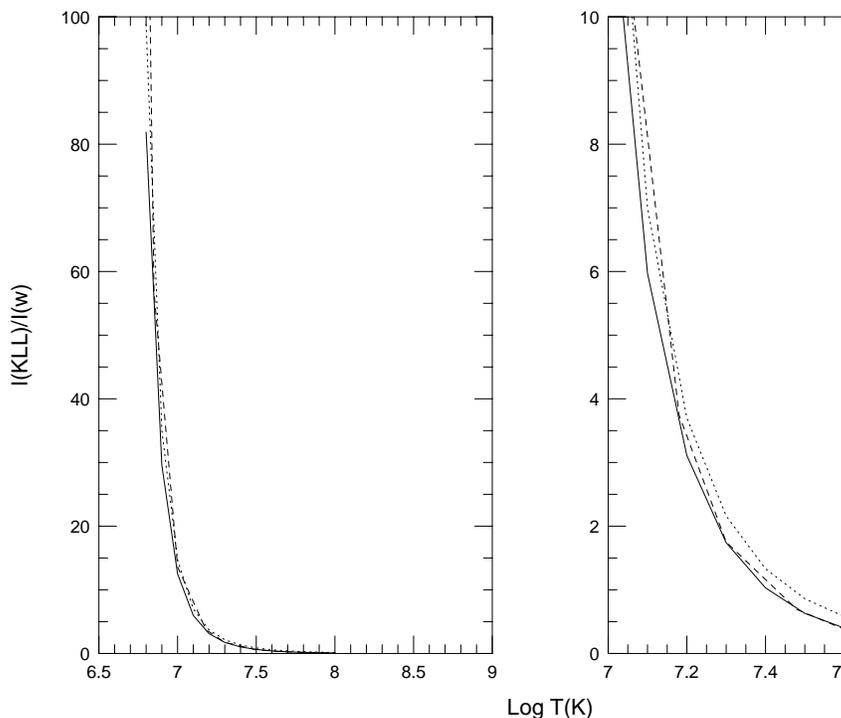


Unified Spectrum of Dielectronic Satellite (DES) Lines



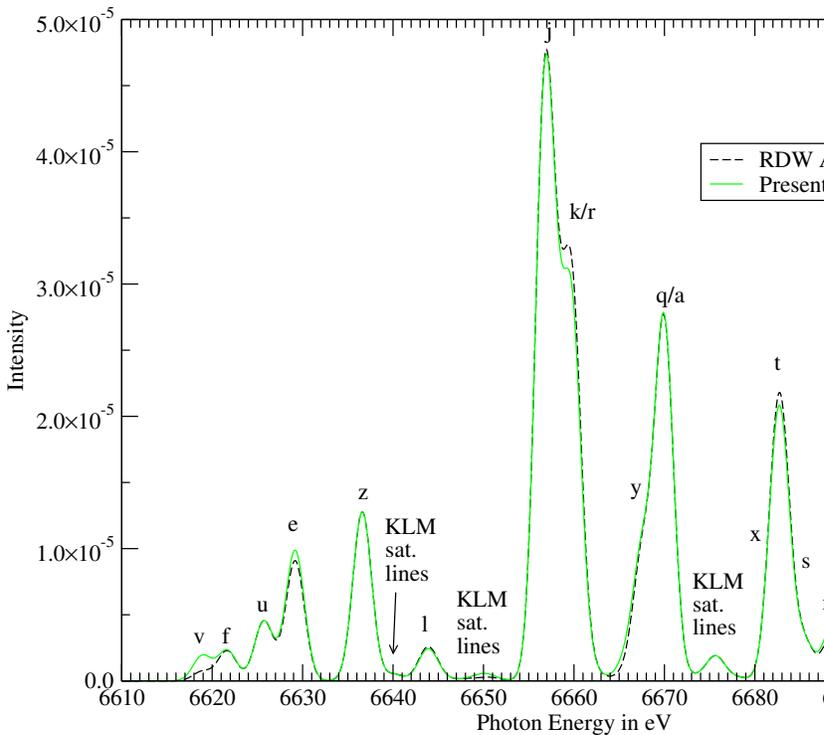
Comparison - DES Intensity ratios $I(KLL)/I(w)$ - Fe XXV (Nahar, Oelgoetz, Pradhan 2007)

solid line – present, dashed line – Bely-Dubau et al. (1982),
dotted line – Vainshtein and Safronova (1978)



- Present unified DES line strengths are in good agreement with those in the individual resonance approximation, BUT some significant differences are found in temperature ranges around $\text{Log T(K)} = 7.2$ where the DES lines are most temperature sensitive, below the temperature of maximum abundance at $\text{Log T(K)} = 7.4$ for Fe XXV in coronal equilibrium.

Comparison of DES spectra of He-like Fe XXV at $T = 1.47 \times 10^7$ K using the relativistic distorted wave (RDW) data from the Los Alamos codes (black) and the present DES strengths computed using the unified recombination scheme (green) (Nahar, Oelgoetz, Pradhan 2007)



- The first application for Fe XXV spectrum shows excellent agreement with that from Los Alamos data.

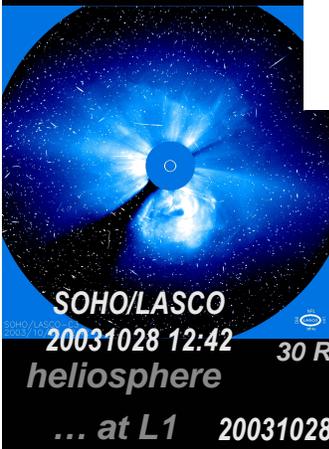
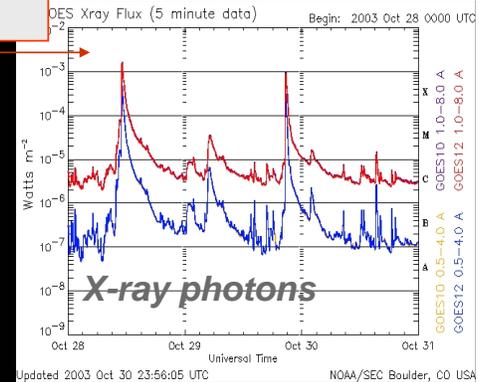
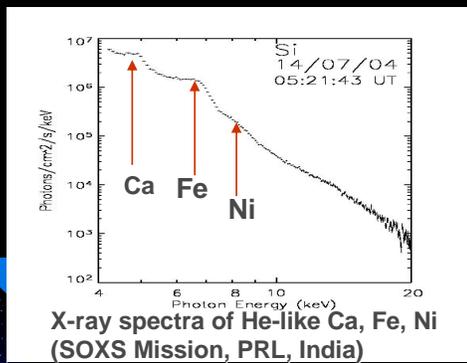
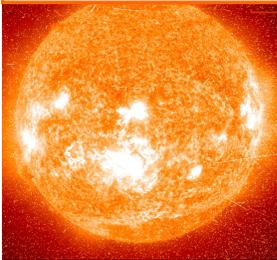
Table: The energies (eV) and strengths of DES lines in KLL resonance complex of (e + Ca XIX) from the unified method.

| Key | Transition | E(eV) | Peak |
|-----|--|--------|-----------|
| o | $1s2s^2(^2S_{1/2}) \rightarrow 1s^22p(^2P_{3/2}^o)$ | 2678 | 4.567E-02 |
| p | $1s2s^2(^2S_{1/2}) \rightarrow 1s^22p(^2P_{1/2}^o)$ | 2678 | 3.225E-02 |
| v | $1s2p^3P^o2s(^4P_{1/2}^o) \rightarrow 1s^22s(^2S_{1/2})$ | 2686 | 5.171E-05 |
| u | $1s2p^3P^o2s(^4P_{3/2}^o) \rightarrow 1s^22s(^2S_{1/2})$ | 2688 | 3.614E-04 |
| r | $1s2p^1P^o2s(^2P_{1/2}^o) \rightarrow 1s^22s(^2S_{1/2})$ | 2716 | 2.143E-01 |
| q | $1s2p^1P^o2s(^2P_{3/2}^o) \rightarrow 1s^22s(^2S_{1/2})$ | 2719 | 3.635E-02 |
| i | $1s2p^2(^4P_{1/2}) \rightarrow 1s^22p(^2P_{1/2}^o)$ | 2726.4 | 3.110E-04 |
| h | $1s2p^2(^4P_{1/2}) \rightarrow 1s^22p(^2P_{3/2}^o)$ | 2726 | 2.134E-05 |
| t | $1s2p^3P^o2s(^2P_{1/2}^o) \rightarrow 1s^22s(^2S_{1/2})$ | 2729 | 3.060E-01 |
| f | $1s2p^2(^4P_{3/2}) \rightarrow 1s^22p(^2P_{3/2}^o)$ | 2729 | 2.616E-04 |
| g | $1s2p^2(^4P_{3/2}) \rightarrow 1s^22p(^2P_{1/2}^o)$ | 2729 | 3.079E-05 |
| s | $1s2p^3P^o2s(^2P_{3/2}^o) \rightarrow 1s^22s(^2S_{1/2})$ | 2731 | 2.033E-01 |
| e | $1s2p^2(^4P_{5/2}) \rightarrow 1s^22p(^2P_{3/2}^o)$ | 2732 | 2.204E-02 |
| k | $1s2p^2(^2D_{3/2}) \rightarrow 1s^22p(^2P_{1/2}^o)$ | 2747 | 9.738E-01 |
| l | $1s2p^2(^2D_{3/2}) \rightarrow 1s^22p(^2P_{3/2}^o)$ | 2749 | 9.559E-01 |
| d | $1s2p^2(^2P_{1/2}) \rightarrow 1s^22p(^2P_{1/2}^o)$ | 2751 | 2.612E-03 |
| c | $1s2p^2(^2P_{1/2}) \rightarrow 1s^22p(^2P_{3/2}^o)$ | 2747 | 7.312E-01 |
| j | $1s2p^2(^2D_{5/2}) \rightarrow 1s^22p(^2P_{3/2}^o)$ | 2749 | 2.192E-01 |
| b | $1s2p^2(^2P_{3/2}) \rightarrow 1s^22p(^2P_{1/2}^o)$ | 2756 | 3.835E-02 |
| a | $1s2p^2(^2P_{3/2}) \rightarrow 1s^22p(^2P_{3/2}^o)$ | 2756 | 2.323E-01 |
| m | $1s2p^2(^2S_{1/2}) \rightarrow 1s^22p(^2P_{3/2}^o)$ | 2774 | 2.320E-01 |
| n | $1s2p^2(^2S_{1/2}) \rightarrow 1s^22p(^2P_{1/2}^o)$ | 2774 | 3.519E-02 |

X-Ray Modeling of Solar Corona and Flares: "Halloween" Solar Storm (Oct 28, 2003)

active region with big sunspot erupts

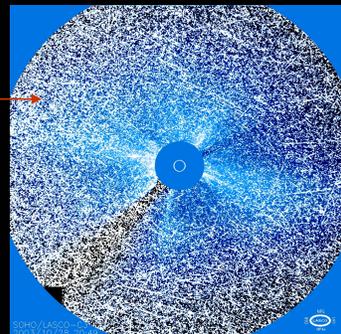
8 minutes later ... X-class Flare observed on the Earth



coronal mass ejection leaves the Sun

8 hours later... particles saturate SOHO/LASCO detector and reach the Earth ("proton shower")

NOAA National Weather Service



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CONCLUSION

1. We present results from ab initio unified method for photoionization and recombination of Ca XVIII, Ca XIX
2. Results includes total and level specific recombination rates and photoionization cross sections of hundreds of bound levels.
3. Total recombination rate coefficients show higher rates at high T compared to previous theoretical results
4. Self-consistent sets of atomic data for photoionization, recombination are obtained and should yield more accurate astrophysical photoionization models